

Aérodynamique Interne: L. Le Penven, M. Buffat, A. Cadiou

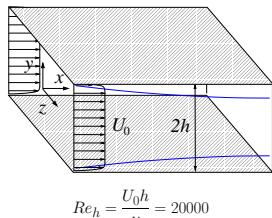
Turbulent transition at the entrance of a plane channel

Motivation

The objective is the study of **by-pass transition** occurring in a plane channel flow, downstream the entry zone, and more particularly the effect of **inlet turbulence** and **wall roughness**. Although the present investigation is currently limited to incompressible flows, one objective is to better understand appearance of cavitation in fuel injector devices (project **NadiaBio** Mov'eo 2008-2011).

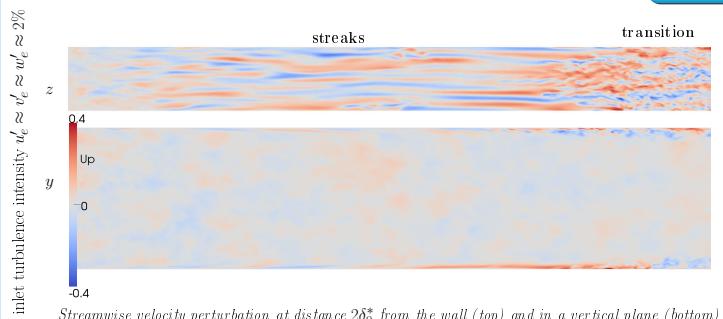
Transition in boundary layers exposed to free-stream turbulent intensity of order $\geq 1\%$ or more can occur without the mediation of viscous Tollmien-Schlichting instability waves (**by-pass transition**). Experiments or numerical simulations show the presence in the boundary layers of large, elongated spanwise modulations of the streamwise velocity called **streamwise streaks**. These disturbances grow in the streamwise direction and are subjected to an instability process leading finally to turbulence breakdown. The physical process explaining emergence of streaks is known as the **lift-up effect**. It is understood as the result of interaction between streamwise vorticity and the boundary layer shear: streamwise vorticity pushes low momentum fluid away from the wall and high momentum fluid towards the wall and the spanwise modulation obtained in this way is stretched downstream by the mean shear. Streaks can be found as solutions of the linearised stability equations (**Orr-Sommerfeld-Squire equations**). However, the effect of viscosity eventually dominates these linear solutions, making the growth of the streaks only transient. Stricto sensu, streaks are thus stable perturbations (sub-critical), but their amplitude can reach such large values (30% of the mean flow) than they escape the linear regime and becomes sensitive to secondary perturbations ([1],[2]).

Flow configuration



$$\delta_0^* = 0.017h : \text{boundary layer displacement at inlet}$$

By-pass transition induced by free-stream turbulence



By-pass transition induced by free-stream turbulence

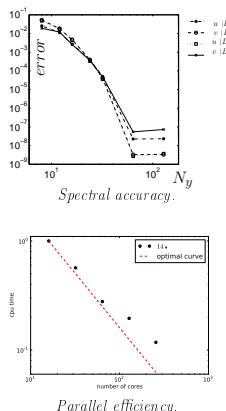
Friction coefficient C_f : transition from Blasius solution to turbulent regime for high enough inlet turbulence intensity.

NadiaSpectral

Using the orthogonal decomposition of the velocity field, an efficient **Galerkin spectral** code has been developed [3] for the analysis of transition in wall-bounded shear flows. In order to use Fourier expansions in the streamwise direction, the solution is forced towards periodicity using the fringe method [4]. The approximation uses **Fourier** expansions in two directions and the **Chebyshev** basis proposed by Moëser et al. [5] in the third direction in order to satisfy the wall boundary conditions. Using Crank-Nicholson/Adams Bashforth time integration, the numerical scheme requires the solution of two 1D sparse linear systems for each Fourier component of the velocity. The method has been implemented in C++ in the **NadiaSpectral** code and parallelized using MPI and OpenMP.

www.ujf-meca.univ-lyon1.fr/~buffat/NadiaSpectral

Numerical method



General result Let be \mathbf{y} a unit vector in \mathbb{R}^3 (the 'vertical' direction) and \mathbf{u} a solenoidal field defined in a bounded domain $\Omega \subset \mathbb{R}^3$, \mathbf{u} can be decomposed into the sum of two solenoidal, $L_2(\Omega)$ -orthogonal fields : $\mathbf{u} = \mathbf{u}_{sq} + \mathbf{u}_{os}$

\mathbf{u}_{sq} is 2D, solenoidal, has zero vertical velocity and is determined by the values of ω , the vertical vorticity of \mathbf{u} ,

\mathbf{u}_{os} is 3D, solenoidal, has zero vertical vorticity and is determined by the values of v , the vertical component of \mathbf{u} and by the values of the normal velocity ω , on the boundary $\partial\Omega$.

For **doubly periodic fields** in the two directions x and z normal to \mathbf{y} : Fourier components can be expressed in function of the components of v and ω . If (α, β) denotes the component of wavevector (α, β) and $k^2 = \alpha^2 + \beta^2$:

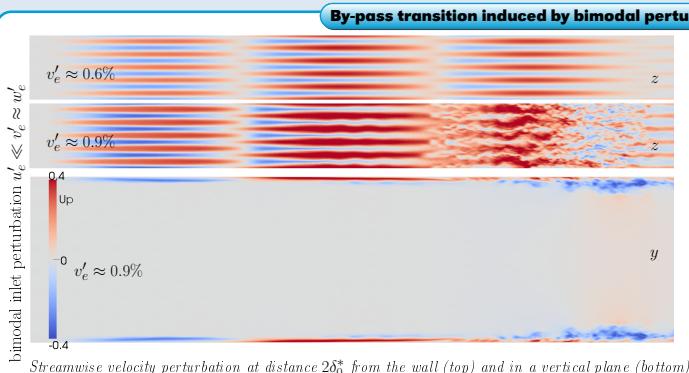
$$\mathbf{u}_{sq}^{\alpha\beta} = \left(-i\frac{\beta}{k^2}\omega^{\alpha\beta}, 0, i\frac{\alpha}{k^2}\omega^{\alpha\beta} \right) ; \mathbf{u}_{os}^{\alpha\beta} = \left(\frac{\alpha}{k^2}\partial_y v^{\alpha\beta}, v^{\alpha\beta}, i\frac{\beta}{k^2}\partial_y v^{\alpha\beta} \right)$$

For plane, parallel, shear flows, this decomposition is related to the **Orr-Sommerfeld and Squire** modal decomposition of the linear stability theory [1]. Using the orthogonal decomposition, the projection of the Navier-Stokes equations linearized around the basis velocity field $(U_b(y), 0, 0)$ reads :

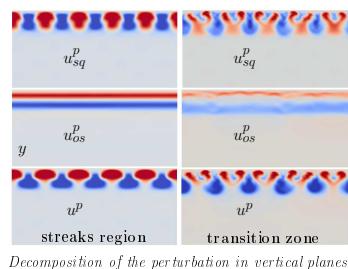
$$\underbrace{\left(\begin{array}{cc} \mathcal{L} & 0 \\ 0 & \mathcal{L} \end{array} \right) \left(\begin{array}{c} \mathbf{u}_{os}^{\alpha\beta} \\ \mathbf{u}_{sq}^{\alpha\beta} \end{array} \right)}_{\text{transport and viscous terms}} + \underbrace{\left(\begin{array}{c} \frac{i\alpha}{k^2}\partial_y (v^{\alpha\beta} U_b') \\ \frac{i\beta}{k^2}(v^{\alpha\beta} U_b') \end{array} \right)}_{\text{interaction with } U_b} = \underbrace{\mathcal{N}(\mathbf{u}_{os}, \mathbf{u}_{sq})}_{\text{non-linear interaction}}$$

The linear part of these two scalar equations is equivalent to the classical Orr-Sommerfeld-Squire system [1].

Orthogonal decomposition of solenoidal fields

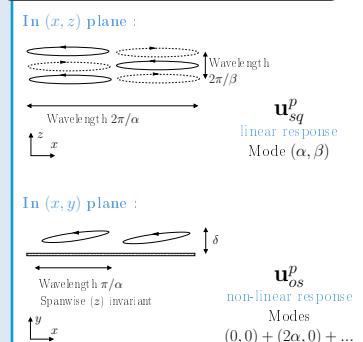


By-pass transition induced by bimodal perturbations : (α, β) (primary) + (α', β') (secondary), $\alpha \ll \beta$



$$u^p = u_{os}^p + u_{sq}^p$$

Non linear structure of periodic streaks



[1] Durbin P. & Wu X., *Transition Beneath Vortical Disturbances*. Annu. Rev. Fluid Mech., 2007

[2] Schmid P. J. & Henningson D. S., *Stability and Transition in Shear Flows*. Springer-Verlag, 2001

[3] Buffat M., Le Penven L., & Cadiou A., *An efficient spectral method based on an orthogonal decomposition of the velocity for transition analysis in wall bounded flow*. Submitted to Computers & Fluids

[4] Bertolotti F. P., Herbert T. & Spalart P. R., *Linear and nonlinear stability of the Blasius boundary layer*. J. Fluid Mech., 1992

[5] Moser R. D., Moin P. & Leonard A., *A spectral numerical method for the Navier-Stokes equations with applications to Taylor-Couette flow*. J. Comp. Phys., 1983

References