# Coupling methods for the convergence rate of Markov processes Joint work with Nicolas Fournier (SU)

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### Outline

#### Coupling and Markovian processes

Motivations Total variation distance and coupling A concrete example: random-to-top shuffling

#### Application to collisionless kinetic theory

Context and main result Stochastic formulation Coupling from the stochastic formulation

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### Motivations

Consider a process satisfying the strong Markov property  $(X_t)_{t \in T}$ with  $T = \mathbb{N}$  (Markov chain) or  $T = \mathbb{R}_+$  (Markov process)

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Assume  $(X_t)_{t \in T}$  takes value in a complete metric space E, endowed with a  $\sigma$ -algebra  $\mathcal{E}$ .

### Motivations

Consider a process satisfying the strong Markov property  $(X_t)_{t \in T}$ with  $T = \mathbb{N}$  (Markov chain) or  $T = \mathbb{R}_+$  (Markov process)

Assume  $(X_t)_{t \in T}$  takes value in a complete metric space E, endowed with a  $\sigma$ -algebra  $\mathcal{E}$ .

Denote  $\mu_0$  for the distribution of  $X_0$ , and assume that  $\mathcal{L}(X_t) \to \pi$ ,  $\pi$  invariant distribution.

Question for today: how can we compute the rate of this convergence ?

## Motivations II

Typical example: you buy a new deck of cards, hence cards are ordered. How long do you need to shuffle them before the distribution is really uniform ?

This question is strongly related to the study of cutoff phenomenon for Markov chains (Diaconis, Shahshahani, Aldous)

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## Motivations II

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For the general question, lots of techniques for Markov chains: Geometric methods (Dirichlet energy), spectral methods, Meyn-Tweedie theory, Wilson's notions of mixing...

Today: focus on one method which translates in continuous time, the coupling method.

### Total variation distance

Consider S the state space, countable for simplicity (but this is absolutely not a restriction). Let  $\mathcal{M}(S)$  the set of probability measures on S.

#### Definition

Let  $\mu, \nu \in \mathcal{M}(S)$ . We define the total variation distance as

$$\|\mu - \nu\|_{TV} = \sup_{A \subset S} |\mu(A) - \nu(A)|.$$

Idea: you try to approximate  $\mu$  with  $\nu$ , and you evaluate your approximation with the worst possible case.

## Coupling: basic definition

### Definition

Given  $\mu, \nu \in \mathcal{M}(S)$ , a coupling of  $\mu$  and  $\nu$  is a couple (X, Y) of random variables such that  $X \sim \mu$ ,  $Y \sim \nu$ .

### Example (Dummy examples)

- If X ~ μ, Y ~ ν with X independent from Y, (X, Y) is a coupling of μ and ν.
- If  $\mu = \nu$ , letting  $X \sim \mu$  and taking Y = X gives a coupling.

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## Coupling: A more interesting example

#### Example

Consider a sequence  $(B_i)_{i\geq 1}$  of independent events with  $\mathbb{P}(B_i) \geq c > 0$  for all *i*. Set  $N = \inf\{i \geq 1, B_i \text{ is realized}\}$ . Then, coupling gives us a way of saying that there exists  $G \sim \mathcal{G}(c)$  such that " $N \leq G$ ".

Indeed, consider a sequence  $(U_i)_{i\geq 1}$  of i.i.d. uniform random variables on [0, 1] and set

$$\alpha_i = \mathbf{1}_{\{U_i \leq \mathbb{P}(B_i)\}}, \quad \beta_i = \mathbf{1}_{\{U_i \leq c\}}.$$

Clearly  $N \stackrel{\mathcal{L}}{=} \inf\{i \ge 1, \alpha_i = 1\}$  and setting  $G = \inf\{i \ge 1, \beta_i = 1\}$ ,  $G \sim \mathcal{G}(c)$ . Moreover,  $\alpha_i \ge \beta_i$  for all i so " $N \le G$ ".

## Coupling and total variation distance

We have the following

Theorem (Relation between coupling and total variation) Let  $\mu, \nu \in \mathcal{M}(S)$ . For all coupling (X, Y) of  $\mu$  and  $\nu$ ,

$$\|\mu-\nu\|_{TV} \leq \mathbb{P}(X \neq Y).$$

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Moreover, there exists a coupling  $(X_*, Y_*)$  such that we have equality.

### An application: Random-to-top shuffling

Assume we have a deck of cards of size n. We label the cards from 1 to n. Hence the state space is

$$S_n = \{ \text{permutations of } \{1, \ldots, n\} \}.$$

We perform a random-to-top shuffling, i.e. at each step we pick a card in the deck at random and put it on top.

By symmetry, assume that  $X_0 = Id$ . Take another chain  $(Y_n)_{n\geq 0}$  such that  $Y_0 \sim \mathcal{U}(S_n) = \pi$ . Since  $\pi$  is invariant,  $Y_n \sim \pi$  for all  $n \geq 0$ .

## Random-to-top shuffling II

At each step, pick  $i \in \{1, ..., n\}$  at random, find the card i and put in on top in both decks .

This is indeed a random-to-top dynamic. Advantage: when card number i is picked, it remains in the same position in both decks for ever .

To have  $X_n = Y_n$ , we thus just need to have selected all cards at time *n*. This is the "coupon-collector" problem: the time to collect all elements of a collection of size *n* when picking uniformly converges in probability towards  $n \ln(n)$ . Set

 $\tau_n = \inf\{t > 0, \text{ all cards have been selected }\}.$ 

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### Random-to-top shuffling III

We conclude that if  $t = (1 + \epsilon)n \ln(n)$  for some  $\epsilon > 0$ , writing  $\mu_k$  the distribution of  $X_k$ ,

$$\|\mu_t - \pi\|_{TV} \leq \mathbb{P}(X_t \neq Y_t) \leq \mathbb{P}(\tau_n > t) \rightarrow 0,$$

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Key ideas to keep in mind: to analyse the convergence towards equilibrium,

- We considered two instances of the process, one of which was distributed according to π.
- We have exhibited a realisation of the dynamic leading to equality (here, selecting all cards).
- We have estimated the time necessary for this realisation to occur.

## Application to collisionless kinetic theory

Consider a gas enclosed in a vessel (bounded domain)  $D \subset \mathbb{R}^n$  with  $n \in 2, 3$ . We study the phase space, hence the density of the gas depends on  $t \ge 0$ ,  $x \in D$ ,  $v \in \mathbb{R}^n$ .

We focus on Knudsen (rarefied) gas: no collisions between particles. Particle in position (x, v) in the phase space just evolves along v until it collides with the boundary.

The randomness comes from the boundary, where particles are reflected stochastically (diffuse reflection).

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### Free-transport equation with boundary condition

The corresponding PDE is given, for  $f_0 \in L^1(D \times \mathbb{R}^n)$ , with  $n_x$  the unit inward normal vector at  $x \in \partial D$ ,

$$\begin{cases} \partial_t f(t,x,v) + v \cdot \nabla_x f(t,x,v) = 0, & (t,x,v) \in (0,\infty) \times D \times \mathbb{R}^n, \\ f(0,x,v) = f_0(x,v), & x \in D, v \in \mathbb{R}^n, \\ f(t,x,v) = cM(v)Kf(t,x), & t > 0, x \in \partial D, v \cdot n_x > 0, \end{cases}$$

where c is a renormalization constant and the flux Kf(t,x) at the point  $x \in \partial D$  at time t is given by

$$Kf(t,x) = \int_{v' \cdot n_x < 0} f(t,x,v') |v' \cdot n_x| dv'.$$

Typically  $M(v) = \exp(-|v|^2) \frac{1}{(2\pi)^{n/2}}$ . In all cases, M has radial symmetry.

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## Main result

We write  $L^1(D \times \mathbb{R}^n)$  for  $L^1(D \times \mathbb{R}^n, \text{Leb})$ . We know that there exists  $f_{\infty} \propto M(v)$ .

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### Main result

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### Theorem (B., Fournier)

Let  $f_0 \in L^1(D \times \mathbb{R}^n)$ , writing  $f_t$  for the unique solution of the problem at time  $t \ge 0$ , if  $r : \mathbb{R}_+ \to \mathbb{R}_+$  is increasing with  $r(x + y) \le C_1(r(x) + r(y))$  for some  $C_1 > 0$  and such that

$$\int_{D\times\mathbb{R}^n} r\Big(\frac{1}{|v|}\Big) f_0(x,v) dv dx + \int_{D\times\mathbb{R}^n} r\Big(\frac{1}{|v|}\Big) M(v) dv dx < \infty,$$

then for all  $t \ge 0$ , for some constant C > 0,

$$\|f_t-f_\infty\|_{L^1}\leq \frac{C}{r(t)}.$$

#### Remark

The solution is a solution in the weak sense of measures.

### Proof strategy

One can show that if  $f_0 \in L^1(D \times \mathbb{R}^n)$ , for all t > 0,  $f_t \in L^1(D \times \mathbb{R}^n)$ . Then the  $L^1$  distance and the total variation distance are equivalent

$$\|f_t-f_\infty\|_{L^1}\propto \|f_t-f_\infty\|_{TV},$$

where we recall that, for  $\mu, \nu$  two probability measures from  $(E, \mathcal{E})$  to  $(F, \mathcal{F})$ , by the previous theorem,

$$\|\mu-\nu\|_{TV} = \inf_{X \sim \mu, Y \sim \nu} \mathbb{P}(X \neq Y).$$

### Proof strategy II

Idea: Construct  $(X_t, V_t, \tilde{X}_t, \tilde{V}_t)_{t \ge 0}$  such that for all  $t \ge 0$ ,  $(X_t, V_t) \sim f_t$  and  $(\tilde{X}_t, \tilde{V}_t) \sim f_{\infty}$ . Then set

 $\tau = \inf\{t > 0 : (X_{t+s}, V_{t+s})_{s \ge 0} = (\tilde{X}_{t+s}, \tilde{V}_{t+s})_{s \ge 0}\}.$ 

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$$\tau = \inf\{t > 0 : (X_{t+s}, V_{t+s})_{s \ge 0} = (\tilde{X}_{t+s}, \tilde{V}_{t+s})_{s \ge 0}\}.$$

Using Markov's inequality we then get

$$\|f_t - f_\infty\|_{L^1} \lesssim \mathbb{P}\Big((X_t, V_t) 
eq ( ilde{X}_t, ilde{V}_t)\Big) \lesssim \mathbb{P}( au > t) \lesssim rac{\mathbb{E}[r( au)]}{r(t)}.$$

To build  $(X_t, V_t)_{t \ge 0}$ , we introduce two functions:

$$\zeta(x,v) = \begin{cases} \inf\{t > 0, x + tv \in \partial D\}, & x \in D \text{ or } x \in \partial D, v \cdot n_x > 0, \\ 0 & x \in \partial D, v \cdot n_x \le 0, \end{cases}$$

$$q(x,v) = x + \zeta(x,v)v.$$

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• Pick 
$$(X_0, V_0) \sim f_0$$
, set  $T_0 = 0$ ,  $T_1 = T_0 + \zeta(X_0, V_0)$ .

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▶ For  $t \in (T_0, T_1)$ , set  $X_t = X_0 + tV_0$ ,  $V_t = V_0$ .

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$$q(x,v)=x+\zeta(x,v)v.$$

 $R_1$  random variable in  $\mathbb{R}_+$  (new norm),  $\Theta_1$  random variable in  $\left(-\frac{\pi}{2}, \frac{\pi}{2}\right) \times [0, \pi]^{n-2}$  such that, for  $x \in \partial D$ ,

$$R_1\vartheta(x,\Theta_1)\sim M(v)|v\cdot n_x|\mathbf{1}_{\{v\cdot n_x>0\}}$$

### A key remark

1. If 
$$V \sim M(v)$$
,  
2. or if  $V \sim \int_D f_0(x, v) dx$ ,  
3. or if  $V \sim M(v) | v \cdot n_x | \mathbf{1}_{\{v \cdot n_x > 0\}}$  for some  $x \in \partial D$ ,  
with the hypothesis of Theorem 1,

$$\mathbb{E}\Big[r\Big(rac{d(D)}{\|V\|}\Big)\Big]<\infty,$$

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where d(D) denotes the diameter of D.

With the previous construction, we get two i.i.d. sequences  $(R_i)_{i\geq 1}$ and  $(\Theta_i)_{i\geq 1}$ . At all time t > 0, the couple  $(X_t, V_t)$  has law  $f_t$ .

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The same construction with  $f_{\infty}$  instead of  $f_0$  gives a process  $(\tilde{X}_t, \tilde{V}_t)_{t\geq 0}$  such that, at all t > 0,  $(\tilde{X}_t, \tilde{V}_t) \sim f_{\infty}$ . We also have two i.i.d. sequences  $(\tilde{R}_i)_{i\geq 1}$ ,  $(\tilde{\Theta}_i)_{i\geq 1}$ .

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The same construction with  $f_{\infty}$  instead of  $f_0$  gives a process  $(\tilde{X}_t, \tilde{V}_t)_{t\geq 0}$  such that, at all t > 0,  $(\tilde{X}_t, \tilde{V}_t) \sim f_{\infty}$ . We also have two i.i.d. sequences  $(\tilde{R}_i)_{i\geq 1}$ ,  $(\tilde{\Theta}_i)_{i\geq 1}$ .

We want to correlate the sequences  $(R_i)_{i\geq 1}$ ,  $(\tilde{R}_i)_{i\geq 1}$ ,  $(\Theta_i)_{i\geq 1}$  and  $(\tilde{\Theta}_i)_{i\geq 1}$  to obtain that  $\mathbb{E}[r(\tau)] < \infty$ .

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We assume D to be strictly convex. We claim that  $\tau \leq T_N$  with  $N \sim \mathcal{G}(c)$  for some c > 0. Indeed, at some time  $T_i$ ,  $i \geq 1$ , we have two possible situations:

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if  $\|\tilde{V}_{T_i}\| < 1$ , we choose the corresponding R,  $\tilde{R}$ ,  $\Theta$ ,  $\tilde{\Theta}$  independently until the data are independent of those at time  $T_i$ . Then we test again.

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if  $\|\tilde{V}_{T_i}\| > 1$  (which happens after a geometric number of iterations of the previous step), we can couple  $(R_i, \Theta_i)$  with  $(\tilde{R}_i, \tilde{\Theta}_i)$  such that, for some  $c_1 > 0$ ,

$$\mathbb{P}\Big((X_{\mathcal{T}_{i+1}},V_{\mathcal{T}_{i+1}})=( ilde{X}_{\mathcal{T}_{i+1}}, ilde{V}_{\mathcal{T}_{i+1}})\Big)\geq c_1.$$

We do so using the maximal coupling from Theorem 1: if  $Z \sim \mu, \tilde{Z} \sim \nu$ , we can correlate Z and  $\tilde{Z}$  such that

$$\mathbb{P}(Z = \tilde{Z}) = \int \mu \wedge \nu.$$

The second step is "successful" after a geometric number of iterations, and we come back to the second step after a geometric number of iterations of the first step (remember the example above !).

A geometric sum of geometric random variables is geometric. The claim follows.

### Conclusion

Once the claim is established we use the previous remark. Upon modifying slightly the rate r, we might use Hölder's inequality: for some  $\epsilon > 0$ ,

$$\mathbb{E}[r(T_N)] \lesssim \sum_{i=1}^{\infty} \sum_{j=1}^{i-1} \mathbb{E}[r(T_{j+1} - T_j)\mathbf{1}_{\{N=i\}}] \ \lesssim \sum_{i=1}^{\infty} \sum_{j=1}^{i-1} \mathbb{E}\Big[r\Big(rac{d(D)}{\|V_{T_j}\|}\Big)^{1+\epsilon}\Big]^{rac{1}{1+\epsilon}} \mathbb{P}(N=i)^{rac{\epsilon}{1+\epsilon}} \ \lesssim \sum_{i=1}^{\infty} i(1-c)^{rac{i\epsilon}{1+\epsilon}} < \infty.$$

For the physically relevant case  $M(v) = \frac{e^{-||v||^2}}{(2\pi)^{n/2}}$ , we find the optimal rate  $\frac{1}{t^{n-1}}$ .

### Extensions

To treat the case where D is not strictly convex, we use a result from Evans (2001): for every  $C^1$  domain there exists  $N \in \mathbb{N}^*$  so that two points at the boundary can be joined in at most Ncollisions with the boundary.

One possible extension is to add specular (deterministic) reflection at the boundary for a fraction of the particles. This changes the construction of the stochastic process, but the proof strategy remains the same.

## Thank you

Thank you for your attention !

