# **Convergence towards the steady state of a collisionless gas with Cercignani-Lampis** boundary condition.

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#### **Problem and main result**

We consider a  $C^2$  bounded domain D in  $\mathbb{R}^n$  for n = 3,  $\partial D$  its boundary. We write  $n_x$  for the unit outward normal vector at  $x \in \partial D$ . Inspired by physical models [KS20] we study the kinetic free-transport equation

> $\partial_t f + v \cdot \nabla_x f = 0 \quad (x, v) \in D \times \mathbb{R}^n,$ (1)

We can thus associate to the problem a semigroup of operators  $(S_t)_{t>0}$  satisfying mass conservation. Given  $f_0 \in L^1(D \times \mathbb{R}^n)$  an initial datum,  $t \ge 0$   $S_t f_0 = f(t, x, v)$  is the solution to (1, 2, 5) at time t. From now on we consider the difference between two solutions, hence  $f_0$ is of mass 0.

# **A Subgeometric Lyapunov Inequality**

We introduce the function



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completed with a boundary condition. We consider the **Cercignani-Lampis** boundary condition, a generalisation of the diffuse boundary condition treated in [Ber20]. Setting  $\partial_{\pm}G := \{(x,v) \in \partial D \times \mathbb{R}^n, \pm (v \cdot n_x) > 0\}, \text{ for all } (x,v) \in \partial_{-}G, t \ge 0,$ 

$$f(t,x,v) = \int_{\{v' \cdot n_x > 0\}} f(t,x,v') |v' \cdot n_x| R(v' \to v;x) dv',$$
(2)

where  $R(u \rightarrow v; x)$  is given, for  $(x, u) \in \partial_+ G$ ,  $v \in \mathbb{R}^n$ , with  $v \cdot n_x < 0$ , by

$$R(u \to v; x) := \frac{1}{\theta(x)r_{\perp}} \frac{1}{(2\pi\theta(x)r_{\parallel}(2-r_{\parallel}))^{\frac{d-1}{2}}} \exp\left(-\frac{|v_{\perp}|^{2}}{2\theta(x)r_{\perp}} - \frac{(1-r_{\perp})|u_{\perp}|^{2}}{2\theta(x)r_{\perp}}\right) \times \exp\left(-\frac{|v_{\parallel} - (1-r_{\parallel})u_{\parallel}|^{2}}{2\theta(x)r_{\parallel}(2-r_{\parallel})}\right) I_{0}\left(\frac{(1-r_{\perp})^{\frac{1}{2}}u_{\perp} \cdot v_{\perp}}{\theta(x)r_{\perp}}\right), \quad (3)$$

where  $v_{\perp} = (v \cdot n_x)n_x$ ,  $v_{\parallel} = v - v_{\perp}$  and the corresponding definitions for  $u_{\perp}, u_{\parallel}$ . Here,  $I_0$  is the modified Bessel function given, for all  $y \in \mathbb{R}$ , by

$$I_0(y) := \frac{1}{\pi} \int_0^{\pi} \exp\left(y\cos\phi\right) d\phi,\tag{4}$$

 $\theta: \partial D \to (0,\infty)$  is the temperature at each point  $x \in \partial D$ , and we have two accommodation coefficients,  $r_{\perp} \in (0,1)$  and  $r_{\parallel} \in (0,2)$ , for, respectively, the normal and the tangential velocities.

 $\sigma(x,v) = \begin{cases} \inf\{s > 0 : x + sv \in \partial D\} \text{ if } (x,v) \in (D \times \mathbb{R}^n) \cup \partial_-G, \\ 0 \text{ otherwise }. \end{cases}$ 

From [Voi80], this function satisfies  $v \cdot \nabla_x \sigma(x, v) = -1$  a.e. in  $D \times \mathbb{R}^n$ , and behaves as  $\frac{1}{|v|}$  when  $|v| \to 0$ . We consider, for  $i \ge 0$ , the weights  $m_i = (1 + \sigma(x, v) + \sqrt{|v|})^i$  and  $||g||_{m_i} = ||gm_i||_{L^1}$  for all  $g \in L^1(D \times \mathbb{R}^n)$ . One can show that

$$\|S_T f\|_{m_i} + \kappa \int_0^T \|S_s f\|_{m_{i-1}} ds \le \|f\|_{m_i} + b(1+T)\|f\|_{L^1}, \tag{6}$$

for  $1 \le i \le (n+1)$ , for all T > 0 and with some constants  $b, \kappa > 0$ . The upper bound on i comes from the necessary condition in the proof that for all  $(x, u) \in \partial_+ G$ ,

$$m_i(x,v)R(u \to v;x)|v \cdot n_x|dv < \infty.$$

Deriving this Lyapunov inequality is the delicate part of the adaptation of the proof from the diffuse boundary condition. Indeed, the boundary operator is no longer weakly compact (we retain some information from the incoming velocity). On the other hand, if v is the incoming velocity, and v' the outcoming one,  $|v'| \le |v|$  on average.

#### **Doeblin-Harris Condition**

Using the nice forms of the characteristics for the problem (1)-(2)-(5), we can show that, for all  $\rho > 0$  there exists  $T(\rho) > 0$  satisfying for some measure  $\nu \neq 0$ ,

 $S_{\mathcal{T}(x)} f > \nu \qquad \qquad f \, dx \, dv, \qquad \forall f \in L^1(D \times \mathbb{R}^n), f > 0.$ 



Figure 1: Diffuse reflection at the boundary, corresponding to  $r_{\perp} = r_{\parallel} = 1$ . Possible outcoming velocities in blue.

This problem corresponds to the behavior of a collisionless gas (Knudsen) enclosed in a vessel. Starting from  $f_0 \in L^1(D \times \mathbb{R}^n)$ , the solution to the problem (1), (2) such that

$$f(0, x, v) = f_0(x, v) \text{ a.e. in } D \times \mathbb{R}^n,$$
(5)

exists and is unique in the  $L^1$  sense. We are interested here in the long-time behavior of this solution: does it converges towards some steady state? If so, what is the rate of this convergence? We prove that

1. there exists indeed a steady state towards which the solution converges in  $L^1(D \times \mathbb{R}^n)$  as  $t \to \infty$ ;

2. the rate of this convergence is  $\frac{1}{(t+1)^{n-}}$  for any reasonable  $f_0$ . This extends the results known for the case  $r_{\parallel} = r_{\perp} = 1$ , and gives another good example of "weak (hypo)dissipativity".



$$\sum_{i=1}^{n} (\rho) J = \nu \int_{\{m_1(x,v) \le \rho\}} J \text{ wave, } v J \in L (L \land \mathbb{R}^n), J \le 0.$$
 (1)

### **Sketch of proof in dimension 3**

We combine (7) and (6) with the method tailored in [CM21]. We derive the following alternative:

$$||f||_{m_2} \le A ||f||_{L^1}$$
, or  $||f||_{m_2} > A ||f||_{L^1}$ ,

for some well-chosen A. In both case, for  $\|.\|_{\alpha,\beta} = \|.\|_{L^1} + \beta \|.\|_{m_{3-}} + \alpha \|f\|_{m_2}$ , we prove that for some  $\alpha > 0, \beta > 0$ ,

$$\|S_T f\|_{\alpha,\beta} \le \|f\|_{\alpha,\beta},$$

from which we conclude  $||S_T f||_{m_{3-}} \le M_3 ||f||_{m_{3-}}$  for all T > 0, some  $M_3 > 0$ . Consider the norm  $\|.\|_{\alpha,\beta,1} = \|.\|_{L^1} + \beta \|.\|_{m_1} + \alpha \|.\|_{m_0}$ . With a similar computation, for some Z constant, using

$$m_1 \leq \lambda m_0 + \epsilon_\lambda m_{3-},$$

with  $\epsilon_{\lambda} = \frac{1}{\lambda^{n-1}}$ , we obtain

$$Z\|S_T f\|_{\alpha,\beta,1} \le \|f\|_{\alpha,\beta,1} + \alpha \frac{\epsilon_{\lambda}}{\lambda} \|S_T f\|_{m_{3-1}}$$

Iterating this result, we conclude that for all  $t \ge 0$ ,

$$\|S_t f\|_{\alpha,\beta,1} \lesssim \left(\frac{1}{(t+1)^{(n-1)-}}\right) \|f\|_{m_{3-1}}$$

reinjecting this in

 $||S_T f||_{\alpha,\beta,1} + 2\alpha ||S_T f||_{m_0} \le ||f||_{\alpha,\beta,1},$ 



Figure 2: The Cercignani Lampis boundary condition in 2D with  $\theta = 1$  and  $r_{\perp} = r_{\parallel} = 0.5$  (left),  $r_{\perp} = r_{\parallel} = 0.1$ (center),  $r_{\perp} = r_{\parallel} = 1/30$  (right). Figures are taken from Chen [Che20].

#### Well-posedness, mass conservation and semigroup

The problem (1)-(2)-(5) satisfies positivity, a result due to the very simple form of characteristics and to the fact that the boundary operator is itself positive. The trace of an  $L^1$  solution f to (1) exists at the boundary, we write it  $\gamma f$ . The mass conservation is easily obtained using the normalization property of R: for all  $(x, u) \in \partial_+ G$ ,

$$\int_{\{v \cdot n_x < 0\}} R(u \to v; x) |v \cdot n_x| dv = 1$$

and iterating, we gain one more exponent to conclude that for all t > 0,



## References

- [Ber20] A. Bernou. A Semigroup Approach to the Convergence Rate of a Collisionless Gas. Kinetic & Related Models, 13(6):1071–1106, 2020.
- [Ber21] A. Bernou. Convergence Towards the Steady State of a Collisionless Gas With Cercignani-Lampis Boundary Condition. Preprint, 2021.
- [Che20] H. Chen. Cercignani-Lampis Boundary in the Boltzmann Theory. Kinetic & Related Models, 13(3):549–597, 2020.
- [CM21] J. A. Cañizo and S. Mischler. Harris-type Results on Geometric and Subgeometric Convergence to Equilibrium for Stochastic Semigroups. Preprint, 2021.
- [KS20] D. Kalempa and F. Sharipov. Drag and Thermophoresis on a Sphere in a Rarefied Gas based on the Cercignani-Lampis Model of Gas-Surface Interaction. Journal of Fluid Mechanics, 900, August 2020.
- [Voi80] J. Voigt. Functional Analytic Treatment of the Initial Boundary Value Problem for Collisionless Gases, January 1980. Habilitationsschrift.